Constraining neutrino mass from galaxy clustering measurements

Matteo Zennaro DiPC, Donostia-San Sebastián

with Julien Bel, Carmelita Carbone, Raúl Angulo, Luigi Guzzo, Jason Dossett





Cosmological neutrinos

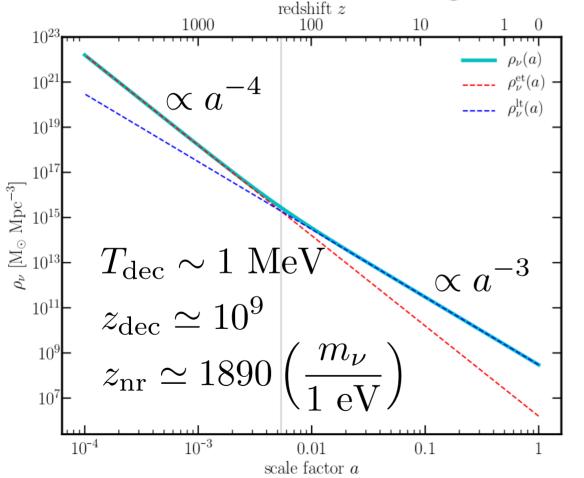
• Oscillation experiments (95%): $M_v = \Sigma m_v > 0.06 \text{ eV}$

Gonzales-Garcia et al (2014), Forero et al (2014), Esteban et al (2017)

- β -decay experiments (95%): m(v_e) < 2.2 eV Kraus et al (2005)
- In cosmology: hot dark matter accounting for a fraction of total dark matter

 $M_v < 0.49 \text{ eV}$ (Planck collaboration, 2015), $M_v < 0.22 \text{ eV}$ (Pellejero-Ibanez et al, 2016), $M_v < 0.12 \text{ eV}$ (Palanque-Delabrouille et al, 2015) and many many more...

Cosmological neutrinos



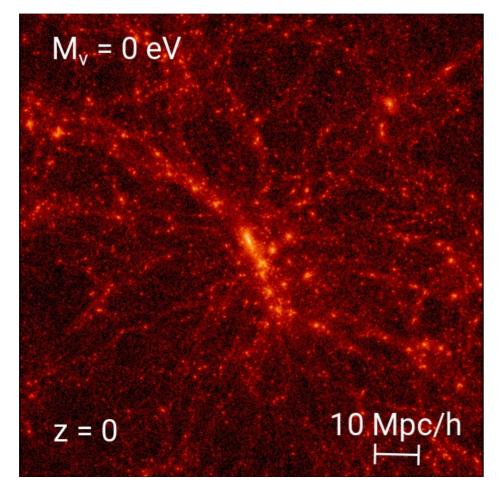
- Light (< 1 eV)
- Weakly interacting
- Fermi-Dirac distribution
- Let's consider:
 - Total mass M
 - Degenerate case $M_v = 3 m_v$

Growth of matter overdensities

$$\frac{\partial^2 \delta_i}{\partial t^2} + 2H \frac{\partial \delta_i}{\partial t} = \frac{c_s^2 \nabla^2 \delta_i}{a^2} + 4\pi G \bar{\rho} \delta_{\text{tot}}$$

$$k_{\rm FS} = \sqrt{\frac{4\pi G\bar{\rho}a^2}{c_s^2}} \simeq 0.91 \frac{\sqrt{\Omega_m(1+z) + \Omega_\Lambda}}{(1+z)^2} \left(\frac{m_\nu}{1\,{\rm eV}}\right) h\,{\rm Mpc}^{-1}$$

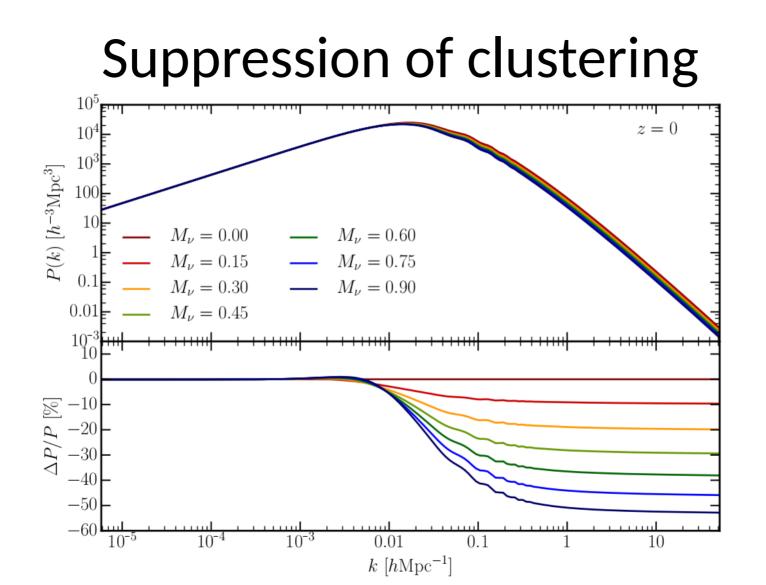
Suppression of clustering



Poisson's equation

$$\nabla^2 \varphi = 4\pi G \bar{\rho} a^2 \left[(1 - \mathfrak{f}_{\nu}) \delta_{\text{cold}} + \mathfrak{f}_{\nu} \delta_{\nu} \right]$$
Effectively \rightarrow 0, for $k > k_{\text{FS}}$

DEMNUni simulations by C. Carbone



Galaxy clustering ratio

Galaxy clustering

- Galaxies are a discrete, biased sampling of the underlying matter field
- If the bias function is local and deterministic

$$\delta_g(x) = F[\delta(x)]$$

• If it is also smooth enough

$$\delta_g(x) = \sum_{i=0}^{\infty} \frac{b_i}{i!} \delta^i(x)$$

Fry & Gaztañaga (1993)

Clustering ratio

Bel & Marinoni (2014)

• 2PCF:

$$\xi_{g,R}(r) = b_1^2 \xi_R(r)$$

• Variance:

 $\sigma_{g,R}^2 = b_1^2 \sigma_R^2$

• CR:
$$\eta_{g,R}(r) \equiv \frac{\xi_{g,R}(r)}{\sigma_{g,R}^2} \equiv \frac{\xi_R(r)}{\sigma_R^2} \equiv \eta_R(r)$$

• z-space (Kaiser):

 $\eta_{g,R}^z(r) \equiv \eta_R(r)$

Clustering ratio: 2nd order

• Hierarchical growth of fluctuations

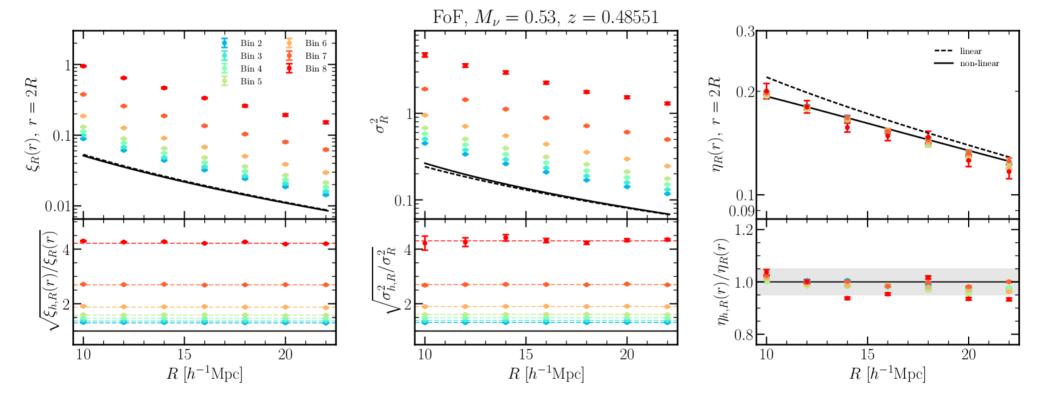
$$\left\langle \delta_R^n \right\rangle_c = S_n \ \sigma_R^{2(n-1)}$$
$$\left\langle \delta_{i,R}^n \delta_{j,R}^m \right\rangle_c = C_{nm} \ \xi_R(r) \ \sigma_R^{2(n+m-2)}$$

• With large enough R, CR is unbiased

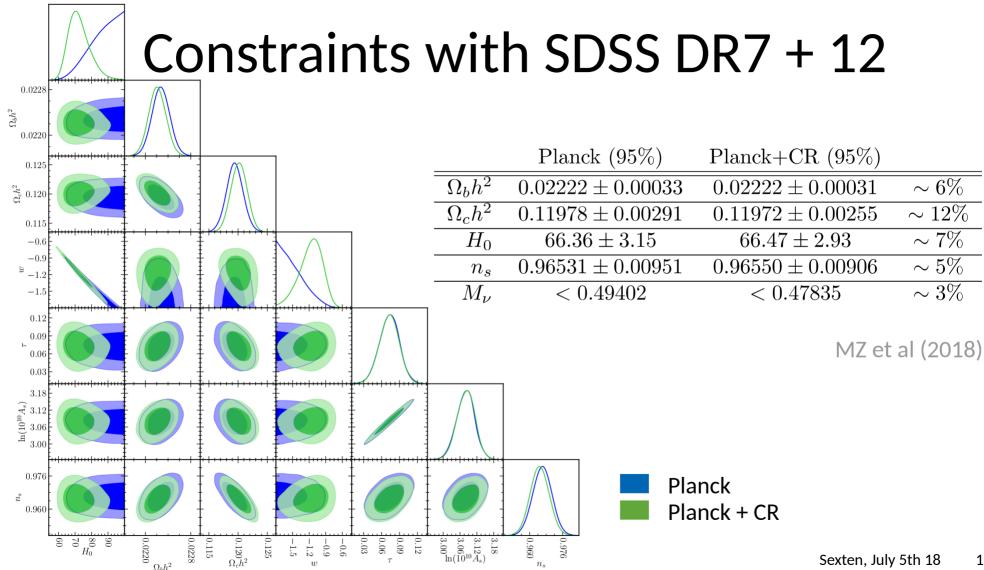
$$\eta_{g,R}(r) \sim \eta_R(r) - \left\{ (S_{3,R} - C_{12,R}) \frac{b_2}{b_1} + \frac{1}{2} \left(\frac{b_2}{b_1} \right)^2 \right\} \xi_R(r) + \frac{1}{2} \left(\frac{b_2}{b_1} \right)^2 \ \eta_R(r) \ \xi_R(r)$$

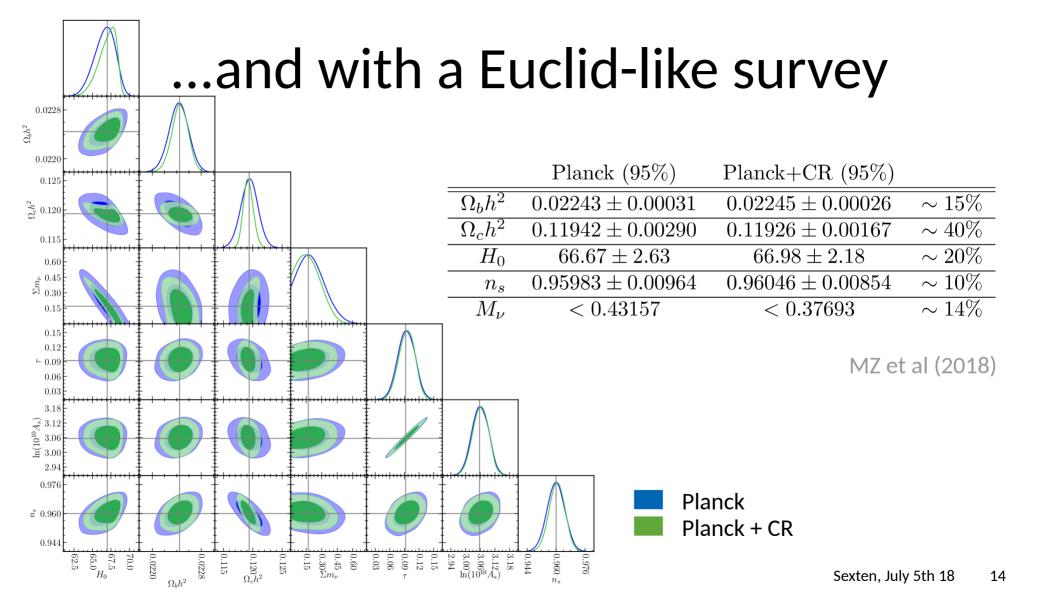
Possible nonlocal contributions only affect b_2 (Bel et al, 2015)

Clustering ratio: unbiased



MZ et al (2018)





Galaxy power spectrum

Galaxy bias: subhalo abundance matching (SHAM)

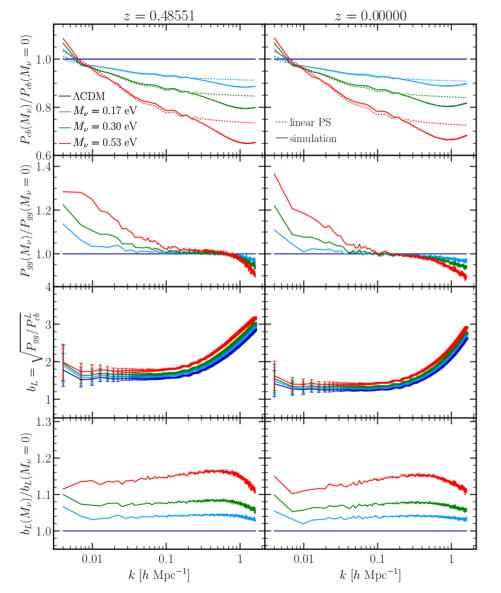
• Matching according to $v_{\max} = \max(\sqrt{GM(\langle r)/r})$

• DEMNUni sims (C. Carbone) $z \simeq \{0.0, 0.5, 1.0, 1.5, 2.0\}$ Carbone et al (2016), Castorina et al (2015), presented today by Andrea Pezzotta

•
$$\bar{n} = \{10^{-3}, 3 \times 10^{-4}, 10^{-4}\} h^3 \text{ Mpc}^{-3}$$

Linear bias model

 $P_g(k) = b_1^2 P(k)$

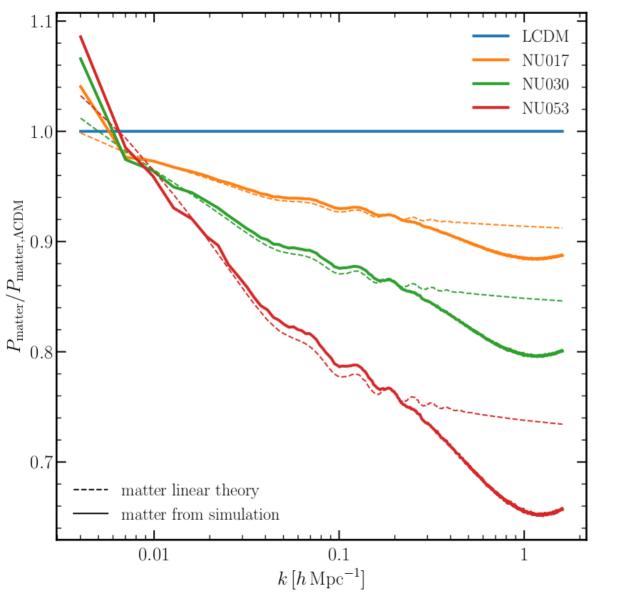


Massive / massless case for cold matter

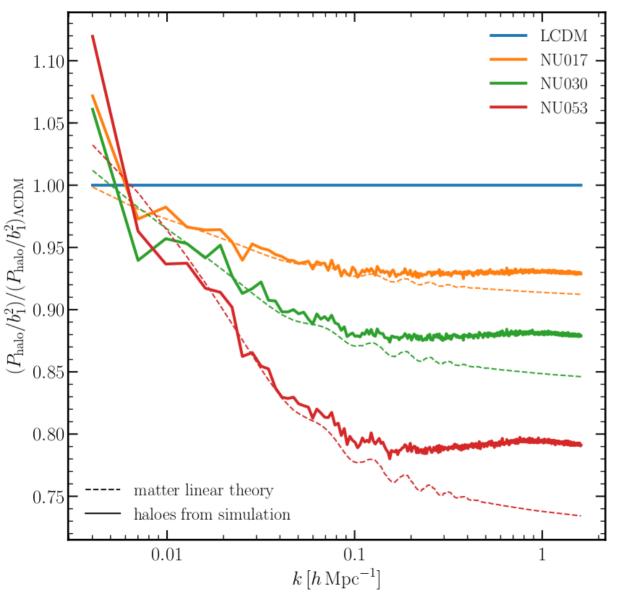
Massive / massless case for galaxies

Linear bias (gaussian errors)

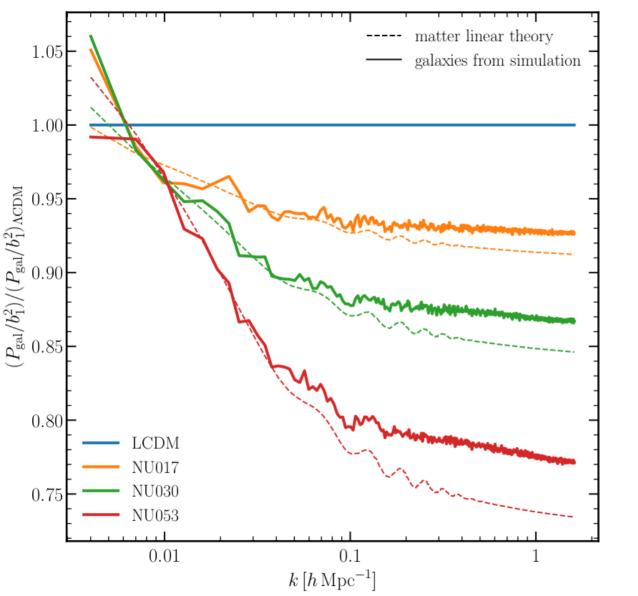
Massive / massless case for the linear bias



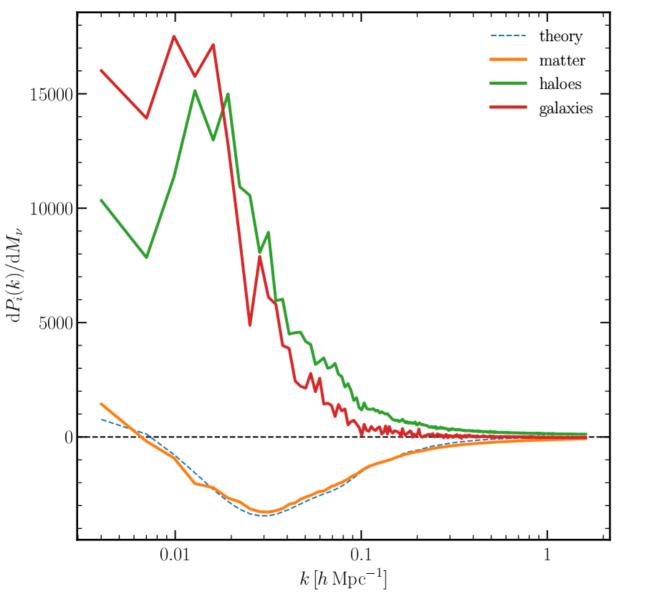




Haloes

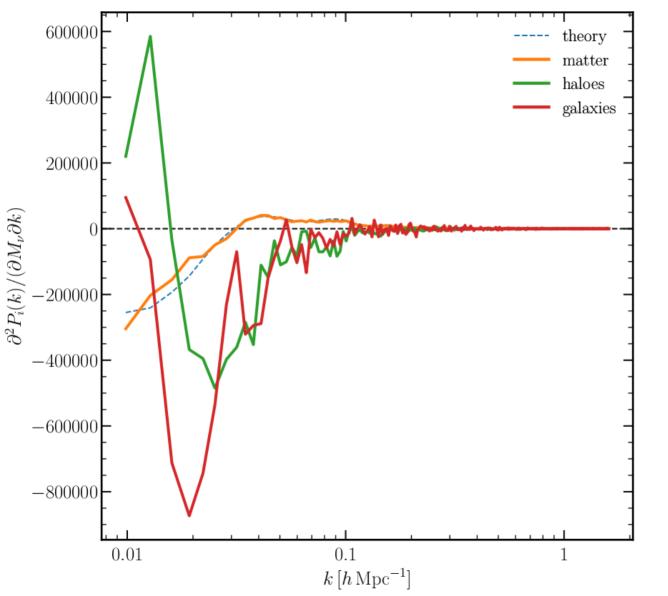


Matter



Derivatives

wrt neutrino mass



Derivatives

wrt neutrino mass and scale

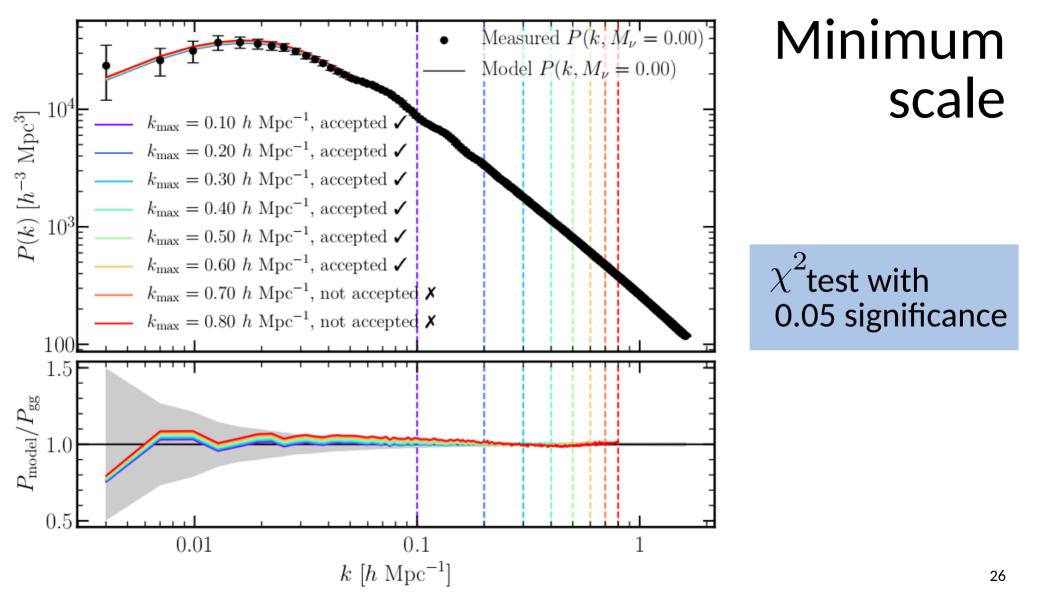
Nonlocal, nonlinear model

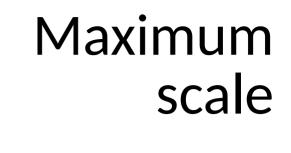
$$\begin{split} P_{\mathrm{g},\delta\delta}(k) &= b_1^2 P_{\delta\delta}(k) + 2b_1 b_2 P_{b_2,\delta}(k) + 2b_1 b_{s^2} P_{b_{s^2},\delta}(k) \\ &+ 2b_1 b_{3\mathrm{nl}} \sigma_3^2(k) P^{\mathrm{lin}}(k) + b_2^2 P_{b_2^2}(k) \\ &+ 2b_2 b_{s^2} P_{b_2 s^2}(k) + b_{s^2}^2 P_{b_{s^2}^2}(k) \end{split} \text{McDonald \& Roy (2009)}$$

Nonlocal, nonlinear: free parameters

- In principle, 4 free parameters $b_1, b_2, b_{s^2}, b_{3nl}$
- If bias is local in lagrangian coordinates
 - Chan et al, 2012: $b_{s^2} = -\frac{4}{7}(b_1 1)$
 - Beutler et al, 2014; Saito et al, 2014: $b_{3nl} = -\frac{32}{315}(b_1 1)$
- Or even relaxing this assumption

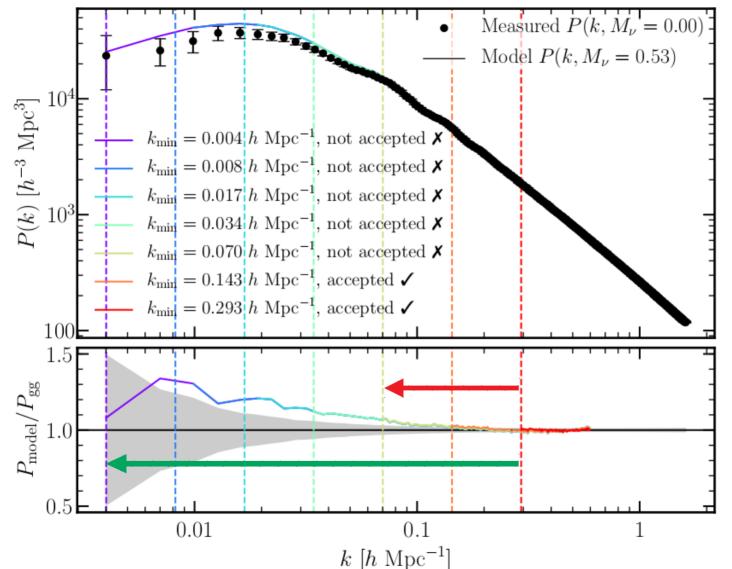
 - Chan et al, 2012: $b_{s^2} = -\frac{4}{7}(b_1 1.43)$ Bel at al, 2015: $b_{s^2} = -\frac{4}{7}(b_1 0.8)$

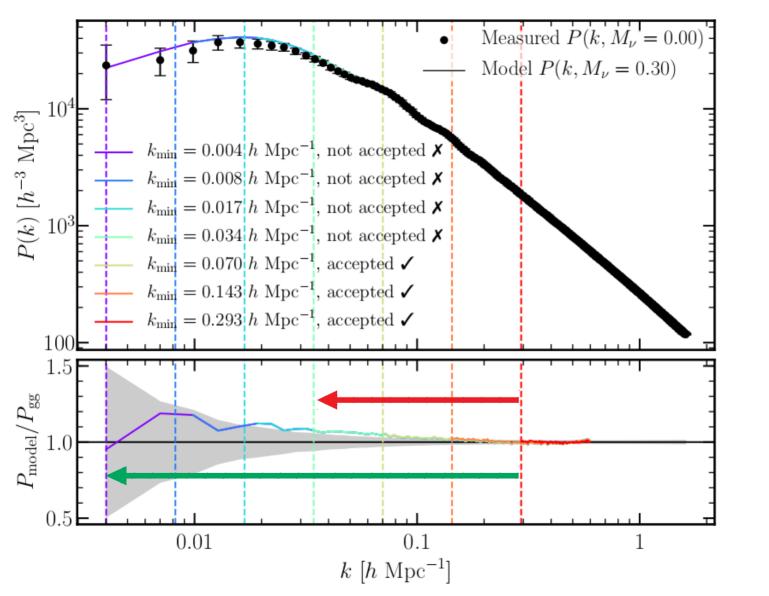




test cannot reject wrong hypothesis

model can distinguish different neutrino mass

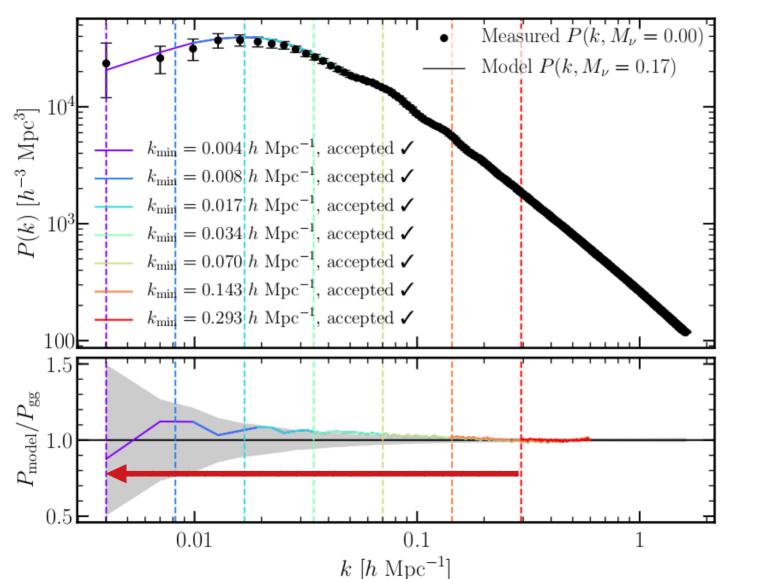




Maximum scale

test cannot reject wrong hypothesis

model can distinguish different neutrino mass



Maximum scale

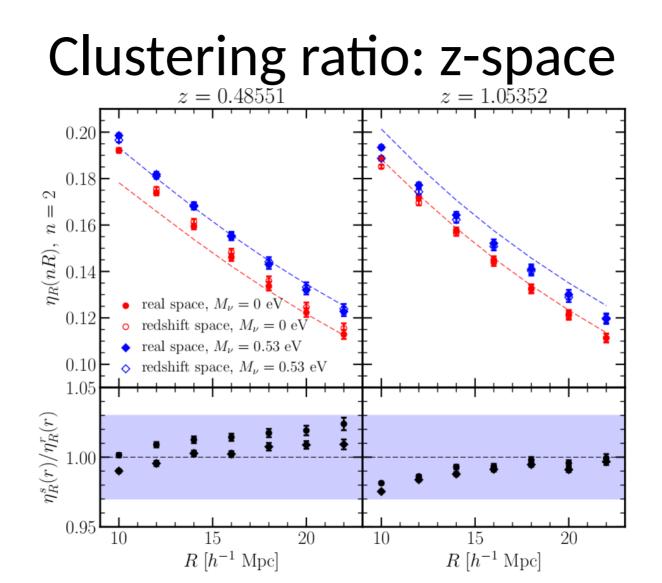
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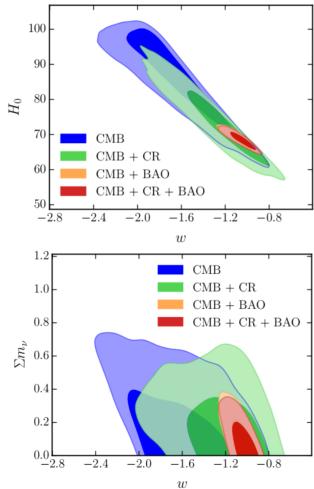
Conclusions

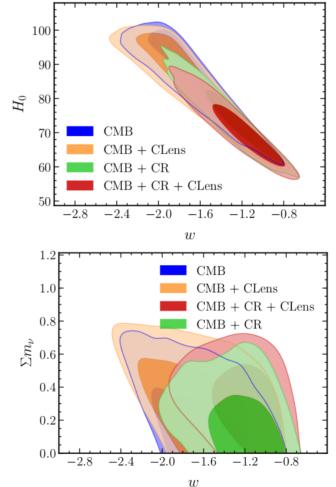
- Extended clustering ratio to cosmologies that include massive neutrinos
- Using SDSS data, 12% improvement on Ω_{cdm} constraints, but not yet competitive for M_{ν}
- Forecasts for Euclid show 40% improvement for Ω_{cdm} and 14% improvement for M_{ν}
- Galaxy power spectrum in the presence of neutrinos with SHAM technique
- Neutrino mass parameter **completely degenerate** if only small scales are included in the fit
- Window of scales where this degeneracy is broken
- MCMC approach to assess the effect on **cosmological parameter constraints**
- **Realistic** galaxies (impact of SHAM flavour, different galaxy populations...)

Backup slides



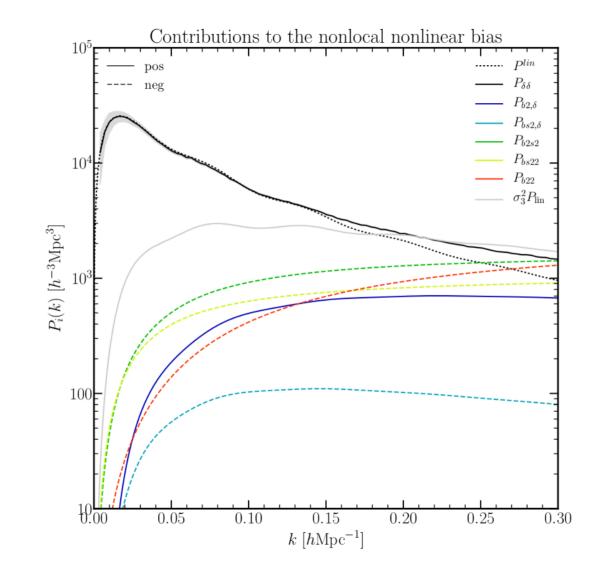
Clustering ratio: CR + BOSS/CLens





Nonlocal, nonlinear: convolution terms

$$\begin{split} P_{b_{2}\delta}(k) &= \int \frac{\mathrm{d}^{3}q}{(2\pi)^{3}} P^{l}(q) P^{l}(|\boldsymbol{k}-\boldsymbol{q}|) \mathcal{F}_{2}^{\mathrm{SPT}}(\boldsymbol{q},\boldsymbol{k}-\boldsymbol{q}), \\ P_{b_{s}^{2}\delta}(k) &= \int \frac{\mathrm{d}^{3}q}{(2\pi)^{3}} P^{l}(q) P^{l}(|\boldsymbol{k}-\boldsymbol{q}|) \mathcal{F}_{2}^{\mathrm{SPT}}(\boldsymbol{q},\boldsymbol{k}-\boldsymbol{q}) S_{2}(\boldsymbol{q},\boldsymbol{k}-\boldsymbol{q}), \\ P_{b_{2}s^{2}}(k) &= -\frac{1}{2} \int \frac{\mathrm{d}^{3}q}{(2\pi)^{3}} P^{l}(q) \left[\frac{2}{3} P^{l}(q) - P^{l}(|\boldsymbol{k}-\boldsymbol{q}|) S_{2}(\boldsymbol{q},\boldsymbol{k}-\boldsymbol{q})\right], \\ P_{b_{s}^{2}}(k) &= -\frac{1}{2} \int \frac{\mathrm{d}^{3}q}{(2\pi)^{3}} P^{l}(q) \left[\frac{4}{9} P^{l}(q) - P^{l}(|\boldsymbol{k}-\boldsymbol{q}|) S_{2}^{2}(\boldsymbol{q},\boldsymbol{k}-\boldsymbol{q})\right], \\ P_{b_{2}^{2}}(k) &= -\frac{1}{2} \int \frac{\mathrm{d}^{3}q}{(2\pi)^{3}} P^{l}(q) \left[P^{l}(q) - P^{l}(|\boldsymbol{k}-\boldsymbol{q}|)\right], \\ \sigma_{3}^{2}(k) &= \int \frac{\mathrm{d}^{3}q}{(2\pi)^{3}} P^{l}(q) \left[\frac{5}{6} + \frac{15}{8} S_{2}(\boldsymbol{q},\boldsymbol{k}-\boldsymbol{q}) S_{2}(-\boldsymbol{q},\boldsymbol{k}) - \frac{5}{4} S_{2}(\boldsymbol{q},\boldsymbol{k}-\boldsymbol{q})\right] \end{split}$$



Sexten, July 5th 18 35